

Superconducting Fluctuations, Pseudogap and Phase Diagram in Cuprates

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Abstract. - We report transport measurements using pulsed magnetic fields to suppress the superconducting fluctuations (SCF) conductivity in a series of $\text{YBa}_2\text{Cu}_3\text{O}_{6+x}$ samples. These experiments allow us altogether to measure the temperature T'_c at which SCF disappear, and the pseudogap temperature T^* . While the latter are consistent with previous determinations of T^* , we find that T'_c is slightly larger than similar data taken by Nernst measurements. A careful investigation near optimal doping shows that T^* becomes smaller than T'_c , which is an unambiguous evidence that the pseudogap cannot be assigned to preformed pairs. Studies of the incidence of disorder on both T'_c and T^* allow us to propose a phase diagram including disorder which explains most observations done in other cuprate families, and to discuss the available knowledge on the pseudogap line in the phase diagram.

Introduction.. – The discovery of High Temperature Superconductivity (HTSC) in the cuprates has been a turning point in the physics of correlated electron systems. Superconductivity (SC) happens indeed in a doped Mott insulator for which nobody would have predicted its occurrence beforehand [1]. While the search of the "pairing glue" remains a very important aim, it has become progressively clear during the last 15 years that it would only be possible to solve this issue after understanding another remarkably robust experimental property of these systems, the pseudogap, which occurs most markedly in underdoped systems. The anomalous behaviour of the electronic properties of these compounds was discovered from a drop of the spin susceptibility $\chi_s(T)$ measured by NMR shifts [2] below a temperature T^* marking the opening of the pseudogap. Specific heat data [3] soon allowed to establish that the density of states is also strongly reduced. The pseudogap has been found generic, that is T^* has a similar variation with doping in bilayer and monolayer [4] systems. Experimentally it was already

clear [2] that T^* decreases so abruptly with increasing doping that it was bound, as shown in Fig.1, either to intersect the SC dome or to merge with the SC line on the overdoped side of the phase diagram (PD). Surprisingly, even for such a simple experimental question, no consensus has been reached so far and the pseudogap line remains nearly evenly distributed between these two possibilities in most recent publications.

More importantly, these two types of PD have been associated with radically different physical [1] explanations of the pseudogap phenomena, either a crossover (or even phase transition) towards a low T correlated state independent or competing with SC [5,6] for Fig.1a, or a simple precursor pairing state, the pseudogap limiting then the T range where SC pairs survive as fluctuating entities above T_c , for Fig.1b.

A strong support for a preformed pair scenario has been initially advocated from the detection of a large Nernst effect [7–9] and of diamagnetism above T_c [10]. Indeed such effects could be associated with superconducting fluctuations (SCF) and/or vortices persisting in the normal state. These experiments introduced a new crossover line which

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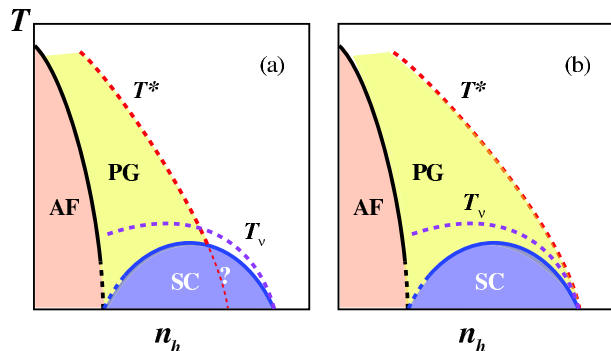


Fig. 1: Possible phase diagrams for the cuprates. The phase transition lines delimiting the AF and the SC states are well established. The two crossover lines lying above T_c signal the opening of the pseudogap and the onset of superconducting fluctuations. Their relative position depends whether a preformed pair scenario or a competing order applies.

one would like to locate in the phase diagrams of Fig.1. As the SCF or vortices require some coherence of the preformed pairs, the initial proposal that the onset of the Nernst effect occurred at $T_\nu \approx T^*/2$, appeared plausible enough to link the pseudogap to the appearance of preformed pairs. However further data [9,11] invalidated this simple relation between T_ν and T^* and established that T_ν does not follow the increase of T^* at low doping but rather drops and appears more related to the T_c dome, especially for systems with reduced intrinsic disorder [9].

This inhibits any conclusion concerning the pseudogap to be taken so far from those experiments. So to discriminate between the PD of Fig.1, one rather needs to find external parameters acting differently on T_c and the two crossover temperatures. Many attempts to locate a crossing of the T^* and T_c lines have been done by exploiting the robustness of T^* to disorder, established early on [12,13]. The reduction of T_c by Zn substitution allowed through analyses of the specific heat [14] or the NMR shift [15] to conclude that, in presence of disorder, any possible T^* value would be below T_c slightly above optimal doping. This has been taken as evidence [14] that the pseudogap line crosses the SC dome and reaches $T = 0$, which would correspond to a quantum critical point (QCP). However this conclusion relies on the *assumption* that, even if T^* merges with the onset of superconductivity, impurities would still suppress T_c and not the pseudogap [15]. But so far no *reliable experiment has ever established the occurrence of T^* values within the SC dome* [13].

The main experimental difficulty remains to distinguish signatures of the SCF and the pseudogap near optimal doping, both for macroscopic probes and for spectral studies by superconducting tunneling (STM) [16] or Angle Resolved Photoemission (ARPES) [17].

In the present work we therefore use a novel experimental approach we established recently [18], which allowed us to suppress the SCF contribution to the ab plane conductivity in single crystal samples with large pulsed magnetic

fields. We do then delineate the onset T'_c of the SCF, which we find located slightly above T_ν and recover the T variation of the normal state resistivity down to T_c from which T^* is determined independently. We exploit here this possibility to determine both T'_c and T^* crossover lines with a single experiment and do give evidence that they intersect around optimal doping, which prohibits considering the pseudogap as the universal onset of pair formation, and favours the PD of Fig.1a. Our results allow us then to discuss the experimental significance of the pseudogaps detected by the spectroscopic experiments.

Samples and techniques. – Single crystals of $\text{YBa}_2\text{Cu}_3\text{O}_{6+x}$ (YBCO) with various oxygen contents have been grown and electrically contacted as described elsewhere [19]. We have studied two underdoped samples (labelled as UD57 and UD85) with T_c (taken at the midpoint of the resistive transition) of 57.1 and 84.6K corresponding approximatively to oxygen contents of 6.54 and 6.8 respectively, an optimally-doped sample (OPT93.6) with $T_c = 93.6\text{K}$ and a slightly overdoped sample (OD92.7) with $T_c = 92.7\text{K}$. The hole doping n_h in these different samples is estimated from the parabolic relationship between T_c/T_c^{max} and n_h [20].

Let us first recall the method we have proposed [18] to recover in high fields the magnetoresistance associated with the normal state quasiparticles. The transverse classical orbital magnetoresistivity in a metal can be written as

$$\delta\rho/\rho_0 = \frac{\rho(T, H) - \rho(T, 0)}{\rho(T, 0)} = (\omega_c\tau)^2 \quad (1)$$

where $\omega_c = eH/m^*$ is the cyclotron frequency and τ an electronic scattering time, so that, as shown by Harris et al [21], the H^2 dependence expected in the low field limit $\omega_c\tau \ll 1$ is obtained in the cuprates at high T for moderate low fields. We have extended this approach by performing $\rho(T, H)$ measurements, using pulse fields as large as 55 Tesla. In view of the required accuracy, we have eliminated any Hall contribution to the measured voltage by reversing the magnetic field. The H^2 dependence could be seen to persist in fields as large as 55 Tesla in the cuprates as long as $T \gg T_c$. As exemplified in Fig.2 for the UD85 sample, $\delta\rho/\rho_0$ indeed displays an H^2 variation for $T > 130\text{K}$, which indicates that the weak field limit still applies up to 55T.

However large departures with respect to this quadratic behaviour appear when T is lowered towards T_c . As already stated [18], we associate the initial fast increase of $\delta\rho/\rho_0$ with the applied field to the destruction by the applied field of the paraconductivity contribution. In such a case the normal state magnetoresistance is only recovered at fields exceeding a T dependent threshold field H'_c . This experimental approach allows us then to single out the normal state properties and separate the contributions of SCF to the in plane transport.

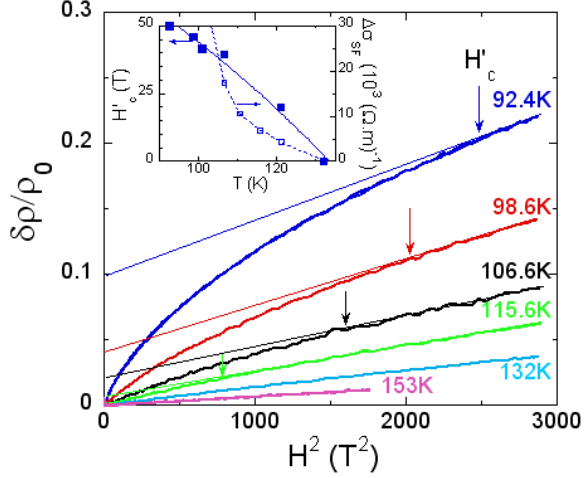


Fig. 2: Field variation of $\delta\rho/\rho_0$ plotted versus H^2 for decreasing temperatures down to $T \approx T_c$ in the UD85 sample. For $T > 130\text{K}$ the magnetoresistance follows an H^2 dependence, but at lower T the contribution of superconducting fluctuations increases the low field conductivity. This occurs below an onset field $H'_c(T)$ (arrow) which is plotted in the inset together with $\Delta\sigma_{SF}(T,0)$. Both quantities become negligible at a temperature $T'_c = 130\text{K}$, which defines the onset of SCF.

Determination of T'_c . – To analyse then the data, it appears quite natural to apply a two fluid model, for which, at a given T , the conductivities due to the normal state quasi-particles $\sigma_n(H)$ and that of the superconducting fluctuations $\Delta\sigma_{SF}(H)$ are additive, that is

$$\rho^{-1}(H) = \sigma(H) = \sigma_n(H) + \Delta\sigma_{SF}(H) \quad (2)$$

At a given T , the limiting high field behaviour $\delta\rho/\rho_0 = \delta\rho_n/\rho_0 + BH^2$ displayed in Fig.2 allows us then to determine

$$\rho_n(H) = \rho_0 + \delta\rho_n + B\rho_0 H^2 = \sigma_n^{-1}(H). \quad (3)$$

Extending such an analysis to the various T values for which the normal state can be recovered with 55Tesla applied field allows us then to determine both $\rho_n(T, H)$ and $\Delta\sigma_{SF}(T, H)$ from Eq.2 and Eq.3. Here we shall concentrate on the conditions for which $\Delta\sigma_{SF}(T, H)$ is fully suppressed, that is for $H > H'_c(T)$. It can be seen in Fig.2 that $\Delta\sigma_{SF}(T, 0)$ vanishes for increasing T at a temperature T'_c which naturally coincides with that for which $H'_c(T)$ vanishes.

This procedure also allows us to extrapolate the zero field resistivity $\rho_n(T, 0) = \rho_0 + \delta\rho_n$ in the absence of SC fluctuations, which is identical to $\rho(T, 0)$ down to T'_c . As an example, the analysis of the UD85 sample data of Fig.2 yields the $\rho_n(T, 0)$ given in Fig.3. Such analyses will be essential to determine hereafter the onset of the pseudogap T^* and therefore will permit direct comparisons of T'_c and T^* .

Let us note that in the pure UD57 samples T'_c has been found larger [18] than the onset of T_v , found for the same

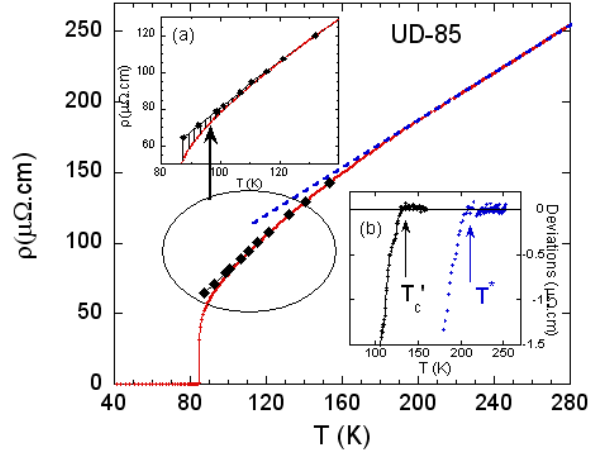


Fig. 3: T variation of the resistivity of the UD85 sample studied. The full square data for $\rho_n(T, 0)$ are deduced here from the pulse field measurements of Fig.2 for $T < T'_c$. The enlargement in inset(a) better exhibits the superconducting fluctuation contribution (hatched area). In inset(b) the resistivity decrease with respect to the high T linear variation due to the opening of the pseudogap is displayed together with that due to SCF.

samples [9], which reflects the fact that the SCF regime is a crossover towards normal state behaviour. The present method is then a quite sensitive experimental approach to study the magnitude of the SCF. The detailed analysis of the T and H dependences of $\Delta\sigma_{SF}$ will be discussed elsewhere [22].

Normal state resistivity versus T , determination of T^* . – In transport measurements, it has often been assumed that the pseudogap can be determined from the downward departure of ρ_{ab} from its linear high T variation [23–25]. In very underdoped cases such as those of UD85 or UD57 one can determine directly T^* without considering the occurrence of SCF at T'_c , as those two crossovers are quite well separated. For instance in UD85, the resistivity deviates from linearity below $T^* \simeq 210\text{K}$, well above $T'_c \simeq 130\text{K}$ as evidenced in Fig.3. This value of T^* agrees well with that ($\simeq 250\text{K}$) obtained from ^{89}Y NMR shift data for similar oxygen contents [12, 26].

However, more care is required for larger dopings, when T^* approaches T_c . If one considers the zero field transport data given in Fig.4a for OPT93.6 and OD92.7, the slopes of the linear T dependences of $\rho(T)$ are quite well defined, and distinct downward deviations are apparent. Those could be associated with the pseudogap T^* if one were to ignore the SCF. However, as can be seen in Fig.4b, the data for $H'_c(T)$ or $\Delta\sigma_{SF}(T, 0)$ give evidence that SCF are the primary source for these downturns, with T'_c values of 135(5) K and 120(5)K respectively (in agreement with diamagnetic measurements for the OPT sample [27]). But one can see on Fig.4a that, when the contribution of SCF is suppressed, the data for $\rho_n(T, 0)$ still displays a down-

ward curvature due to the pseudogap. This is better displayed in Fig.4c, in which the deviation of $\rho_n(T, 0)$ from linearity has been magnified. There one can determine $T^* = 118(5)\text{K}$ for OPT-93.6, while one cannot ascertain any deviation for OD92.7 within experimental accuracy, so that $T^* < 100\text{K}$ in that case. Notice that we consistently defined here T'_c and T^* at temperatures for which deviations of $\Delta\sigma_{SF}(T, 0)$ or $\rho_n(T, 0)$ correspond to the same magnitude for both hole contents (hatched areas in Fig.4b and c).

So, while T'_c is always smaller than T^* for underdoped samples, the present data give evidence that $T^* < T'_c$ for OPT93.6, while for OD92.7 any possible value for T^* would only also occur below T'_c . This establishes that the pseudogap line crosses the onset of SCF for a hole content slightly below optimal doping, as shown in Fig.4d, where the data for T'_c and T^* are summarized versus the equivalent hole contents. This result definitively *prohibits considering that the pseudogap could be the onset of pairing*. We shall discuss its importance after showing the incidence of disorder on these two lines in the phase diagram.

Influence of disorder. – It is well established that disorder strongly depresses T_c while T^* is not affected [12, 13]. This can be directly evidenced from resistivity measurements in electron irradiated crystals. Electron irradiation at low T has been shown to introduce homogeneously distributed point defects without changing the hole doping [28, 29].

For an UD59 single crystal, Fig.5a shows that $d\rho_{ab}/dT$ increases with decreasing T below 300K, which would correspond to $T^* > 300\text{K}$, in agreement with $T^* \simeq 350$ to 400K found from NMR [2, 12, 14] for all samples with oxygen contents between 6.6 and 6.7 at the $T_c \approx 60\text{K}$ plateau. In such samples one cannot use the deviation of $\rho_{ab}(T)$ from linearity to determine T^* , as diffusion of oxygen takes place above room temperature yielding significant variations of the chain contributions to $\rho(T)$ [30]. However it has been shown that the inflexion point in $\rho(T)$ is located at $T_1 \simeq T^*/2$ [24, 25]¹. The fact that this point only slightly increases with increasing defect content is a good confirmation of the NMR results that T^* is not modified by disorder. This is also a different phrasing of the evidence that Matthiessen's rule applies in the normal state of cuprates [28]. On the contrary, we have previously shown from Nernst effect data [9] and high field measurements [18] that disorder depresses T_ν and T'_c , although in much less proportion than T_c .

We have also performed pulse field measurements on an OPT-93.6 sample with T_c reduced down to 70.7K by electron irradiation in order to test the influence of disorder on both T'_c and T^* . As can be seen in Fig.5b, T'_c is found there to decrease markedly down to 95K. There-

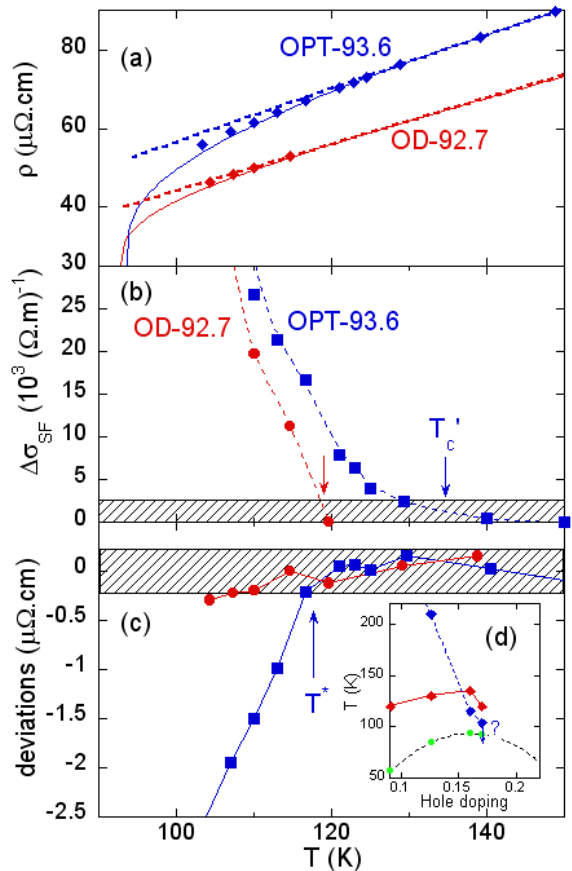


Fig. 4: (a) T variations of the zero field resistivities and $\rho_n(T, 0)$ deduced from high field data for OPT-93.6 and OD92.7 samples. (b) Onset of the SCF contribution to the conductivity allowing to determine T'_c . (c) Deviations from linearity of $\rho_n(T, 0)$ allowing to determine the pseudogap T^* . (d) Phase diagram versus hole doping showing that the crossover lines for T'_c and T^* intersect around optimal doping.

fore, above 95K, the data for $\rho(T)$ is not influenced by SCF for this sample, and the detected deviation from the high T linear behaviour gives $T^* = 120(5)\text{K}$. This value is in good agreement with that obtained above in the pure sample after suppressing the SCF in high fields. So we confirm here that, as previously shown in UD57 sample, defects depress T_c and T'_c near optimal doping, but have a much smaller incidence on T^* , which shows that these two crossover temperatures are associated with independent physical phenomena.

Discussion. – The use of high applied fields has been essential in allowing us to depress the SCF in pure YBCO samples, which permits us to shed some light on the continuing debate concerning the pseudogap and SCF in cuprates. We do demonstrate that over the hole concentration range from underdoped to slightly overdoped, the SCF crossover T'_c occurs at quite similar temperatures, and is affected in the same manner by large applied fields.

¹Ando et al. [31] propose another determination of T^* based on the variation of the second derivative of the resistivity. This leads to much lower values of T^* than those determined by NMR.

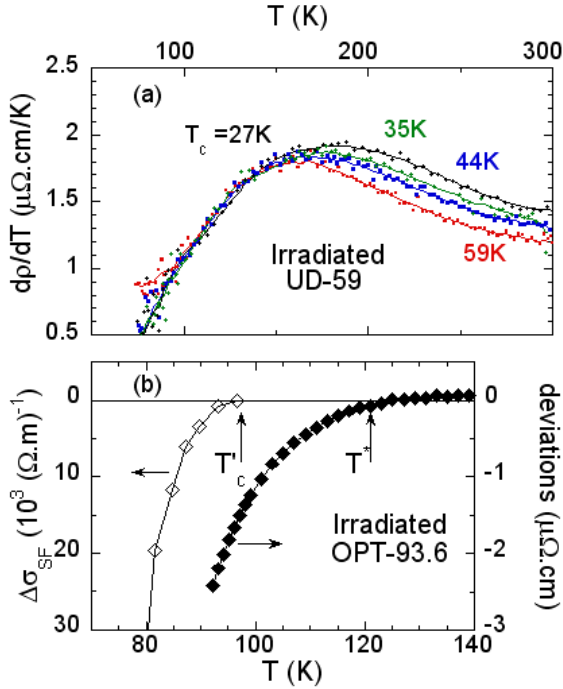


Fig. 5: (a) Variations of $d\rho_{ab}/dT$ versus temperature in the pure UD59 sample and after electron irradiation at low T . The maximum of $d\rho_{ab}/dT$, which roughly correspond to $T^*/2$ does not vary markedly with disorder. (b) Similarly in OPT93.6, pulse field measurements gives evidence that T'_c is depressed to 95K by disorder, while T^* remains nearly unmodified at ≈ 120 K as in the non irradiated sample of Fig. 4.

On the contrary the pseudogap crossover T^* is shown to depend markedly on doping and to cross T'_c at optimal doping, so that the onset of pairing occurs there above T^* . Here we have consistently located T^* and T'_c for different dopings at temperatures corresponding to resistivity deviations of the same magnitude.

Furthermore we confirmed, within the same experimental procedure, that disorder does not affect T^* while it depresses T'_c , so that the clean case phase diagram displayed in the inset of Fig.4c can be extended in a third dimension by introducing a disorder axis in Fig.6, for realistic ranges of disorder. There we display that i) the T_c depression is faster for underdoped samples than for optimal ones, ii) the T^* line does not evolve with disorder, iii) the onset T'_c of SCF decreases at a much slower rate in the underdoped case than for optimal doping.

Consequently the SCF regime appears then much more extended with respect to T_c values in the underdoped samples when a large disorder is introduced. Such fluctuations were even detected in underdoped samples for which T_c was nearly reduced to zero [18]. This PD therefore mimics the situation encountered for Nernst data taken in $\text{La}_{2-x}\text{Sr}_x\text{CuO}_4$, for which the SCF regime was found to be peaked in the underdoped range [8]. This agrees with our former proposal that large intrinsic disorder explains

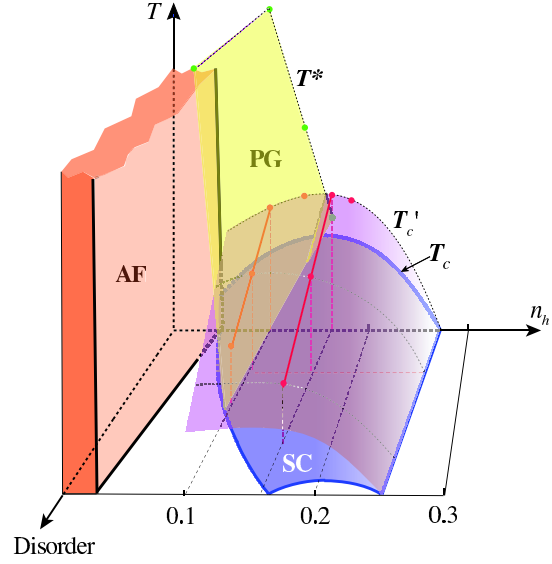


Fig. 6: (color on line)Phase diagram constructed on the data points obtained here, showing the evolution of T'_c the onset of SCF, with doping and disorder. The fact that the pseudogap and the SCF surfaces intersect each other near optimum doping in the clean limit is apparent. These surfaces have been limited to experimental ranges where they have been determined experimentally. As discussed in the text, we did not represent any putative evolution of T^* below T_c , nor of T'_c towards the regimes where static spin glass like magnetic disorder appears. In the overdoped regime, data taken on Tl 2201 indicates that disorder suppresses SC without any anomalous extension of the SCF [28].

the metal to insulator crossover seen in that case. This could be opposed to the observation that YBCO remains in a well defined metallic state at low T for similar hole dopings [32].

The specific behaviours of the SCF and of the pseudogap remain among the most important features of the cuprate physics. It is quite clear that, independently of the actual possible origin of the pseudogap, SCF are present above the 3D T_c in these 2D systems with short coherence length. One would still like to understand whether SCF are purely gaussian amplitude fluctuations of the order parameter, or whether the loss of coherence of the SC state is achieved by phase fluctuations at least in part of the phase diagram. The quantitative analysis of our data on the SCF conductivity will allow us to discuss those points elsewhere [22].

As for the pseudogap phenomenon, anomalies of distinct transport experiments have been detected above T_c and assigned as well to the pseudogap, locating it at much higher T^* values for optimal doping than found here. This has been the case for an increase in the c axis conductivity [33] or for the onset of in plane anisotropy of the Nernst coefficient [34] in YBCO samples. While in plane transport and magnetic susceptibilities have been studied at length in a variety of cuprates for the last twenty years,

these recent experiments probe more complex transport properties, which are hardly understood in other layered oxides. The reproducibility of the detected anomalies in other HTSC, and their insensitivity to disorder have not been established so far, so that their use to locate T^* does not appear to us as secure as the present approach.

Meanwhile STM [16] or ARPES [17] spectroscopic experiments, have detected gaps in the energy spectrum, which are found to persist above T_c in underdoped but also in overdoped cuprates. The excited states detected by STM experiments are easier to analyze in the overdoped cuprates in the absence of pseudogap [35,36]. The d-wave superconducting gap detected at low T does not fully disappear above the 3D T_c , so that the remaining dips in the energy spectrum can only be associated with SCF in this doping range.

For optimally doped samples the SCF may extend below or above the onset of the pseudogap, as shown here. But it is not possible so far to use large enough fields to separate them in the spectroscopic ARPES and STM experiments. This explains why remnants of the SC gap due to SCF could be often taken inadvertently as "pseudogaps".

We have shown that the robustness of the pseudogap to disorder is an important clue to separate its spectral features from those associated with SC. In the STM and ARPES experiments for which surface inhomogeneities are apparently extremely important [37], it seems to us that the large gap detected in underdoped samples ought to be associated with the pseudogap, in view of its robustness and reproducibility. The Fermi arcs (or pockets) are also undoubtedly experimental features associated with the pseudogap state.

The pseudogap being distinct from SCF, one then expects to detect different spectral responses by STM or ARPES for these two phenomena below optimal doping. But for underdoped samples experimental difficulties could have been anticipated from the early experiments on the cuprates. Indeed, NMR experiments have shown for long that the depression of χ_s at low T is so large that it was nearly impossible to detect the further reduction of χ_s expected below T_c . So if the large gap detected in the underdoped systems is indeed associated with the pseudogap, it renders quite difficult the detection of spectral modifications occurring then below T'_c . This explains the debate between one gap [38] and two gap scenarios [39,40] which has been engaged for a while, and for which no consensus has emerged so far.

Finally, if any feature detected above T_c is not always connected with the pseudogap, one may wonder as well if anything detected below T^* is an essential phenomenon or a secondary transition which might be non universal to the cuprate physics and driven by specific system details. If T^* is connected with a competing order which breaks time reversal symmetry [41,42], one still needs to establish the correlation between the various experimental aspects of the pseudogap -loss of spin susceptibility, transfer of spectral weight from low to high energies, Fermi arc

observations,...- and the symmetry breaking.

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